

DFT Analysis of Ca-doped LiFePO₄ Cathode Material for High-Performance Li-ion Batteries

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This study investigates the effects of calcium (Ca) doping on lithium iron phosphate (LiFePO₄) using the first principles method. Despite its stability and safety, its practical use is limited by low electronic conductivity and sluggish lithium ion diffusion. To address this issue, density functional theory (DFT) simulations were used to investigate how Ca doping influences the mechanical, electrical, and electrochemical properties of LiFePO₄. The results demonstrated that Ca doping reduced the energy gap of LiFePO₄ and FePO₄ to 2.164 eV and 1.110 eV, which led to better electron transport and improved electronic conductivity in the cathode material. Additionally, it increased the lithium diffusion coefficient from $1.04 \times 10^{-11} \text{ cm}^2 \text{ s}^{-1}$ to $1.75 \times 10^{-10} \text{ cm}^2 \text{ s}^{-1}$, which improved lithium ion mobility and rate capability. There was a slight expansion in lattice parameters which may facilitate better lithium diffusion pathways. Based on its mechanical properties, Ca-doped LiFePO₄ remained stable and there was increased ductility and isotropy which may help prevent microcrack formation during prolonged electrochemical cycling. Overall, the modification enhanced the performance of the LiFePO₄ cathode material, paving the way for advancements in battery technology.

Keywords: Cathode material; lithium-ion battery; calcium; doping; LiFePO₄

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Since its discovery by Padhi and coworkers of John B. Goodenough in 1996, lithium iron phosphate (LiFePO₄) has emerged as a promising cathode material for rechargeable lithium-ion batteries due to its inherent safety, stability, and environmental friendliness [1]. This olivine LFP (LiFePO₄) has a lithium intercalation potential of 3.5 V with a theoretical capacity of 170 mAhg⁻¹ [2]. Despite these advantages, LFP suffers from low electronic conductivity ($\sim 10^{-9}$ – $10^{-10} \text{ S cm}^{-1}$) and sluggish lithium ion diffusion rates ($\sim 10^{-14} \text{ cm}^2 \text{ s}^{-1}$), which limits its use in high-power applications [3–5]. To improve these properties, various strategies have been explored, including surface carbon coating, nanosizing, and doping with various elements [6–9].

Among these, carbon coating is the most widely used due to its cost-effectiveness and ease of large-scale production. This approach enhances conductivity by encapsulating the LFP surface and forming an electronic conductive network. Generally, the surface modification approach does not interfere with the internal material and has no significant impact on the material's intrinsic problems. LFP's poor conductivity is mainly due to its structure which only allows one

dimensional lithium diffusion, leading to sluggish ion diffusion [10–12]. Despite the strong covalent P–O bond which stabilizes the structure, this characteristic hinders Li⁺ transport into the PO₄³⁻ polyanion, thereby restricting the pathway to a single dimension. Moreover, the absence of a continuous network of Fe–O inside the crystal structure contributes to low electron conductivity [13].

Doping has been explored as an alternative strategy to modify the electronic structure of LFP and enhance its electrochemical performance. Various dopants, including transition metals (Mn, Co, V, Nb, Cr), heteroatoms (F, Cl, S), and alkali/alkaline earth metals (Na, K, Ca), have been investigated to understand the effects of doping on the properties of LiFePO₄ [14–20].

Among the potential dopants, calcium (Ca) has gained attention due to its favourable impact on the structural and electronic properties of the cathode material. Qu and coworkers [21] investigated the effect of Ca doping on the performance of Li₂FeSiO₄/C. They found that Ca doping significantly improved

the electrochemical properties of the Li₂FeSiO₄/C material. With the introduction of Ca atoms, there was a noticeable enhancement in structural stability, which in turn improved the material's electronic conductivity, effectively reduced the charge-transfer resistance, and increased the diffusion coefficient of lithium ions, which may lead to better capacity retention and rate capability. Previous work on Ca doping in LiMnFePO₄ showed that it significantly enhanced battery performance [22]. Ca doping increases electronic conductivity and the lithium ion diffusion coefficient, resulting in improved overall conductivity and faster transport within the material. Moreover, Ca-doped LiMnFePO₄ exhibits superior cycling stability, maintaining higher capacities over more charge-discharge cycles compared to its undoped counterpart. Kugnatahan et. al [23] studied the effects of various doping elements. They found that Ca doped at the Fe site of LFP was thermodynamically stable with a solution energy of -0.42 eV, making Ca the most favourable dopant among alkali-earth metals like Mg, Sr, and Ba, partly due to the ionic radius (1.00 Å) of Ca²⁺ which is close to that of Fe²⁺ (0.78 Å), while its Bader charge (+1.60) is almost equal to that of Fe (+1.55). This suggests that synthesizing Li(Fe,Ca)PO₄ is feasible and could enhance the electrochemical properties of the LFP cathode material.

Despite promising results, systematic computational studies on Ca-doped LiFePO₄ remain limited. This study aims to bridge that gap by using density functional theory (DFT) simulations to analyse how Ca doping affects the structural, electronic, and electrochemical properties of LiFePO₄. This study investigates how Ca incorporation alters the lattice structure, electronic properties, and lithium-ion migration energy of LFP, which can provide new insights into the potential of Ca-doped LFP for high-performance lithium-ion batteries.

METHODOLOGY

In this work, computational studies were carried out based on density functional theory (DFT) implemented in the Cambridge Serial Total Energy Package (CASTEP) [24] computer code available in Material Studio software. The spin-polarized general gradient approximation (GGA) with Perdew-Burke-Ernzerhof (PBE) [25] was used for the exchange-correlation function. The ultrasoft pseudopotential was used to model valence-electron interactions with Li 1s¹2s², O 2s²2p⁴, P 3s²3p³, Fe 3d⁶4s², and Ca 3s²3p⁶4s² considered as valence electron configurations. The effective Hubbard U parameter of U = 4.3 eV was set to the Fe 3d electronic state to correct the famous underestimation of the electronic band gap by conventional DFT, known as the GGA + U method. The plane wave cut-off energy was set to 650 eV and the Monkhorost-Pack scheme k-point of 1×1×1 was employed for the integration of the Brillouin zone. A supercell of 1×2×2 was used, which comprised of

four units of LFP cells (Li₁₆Fe₁₆P₁₆O₆₄) with 112 atoms. One of the Fe atoms was replaced by a Ca atom representing 6.25 % impurities, which was the lowest Ca concentration at the Fe site that could be obtained in this system. The chosen cell size provided a practical balance between representing the doped system accurately and maintaining computational efficiency. While experimental studies typically investigate Ca doping levels in the range of 1–5 % [21, 22, 26], slightly higher concentrations are often employed in theoretical studies due to the computational costs associated with larger supercells required for lower dopant levels. The FePO₄ (FP) structure was constructed by removing all the lithium ions from the cell. The Broyden–Fletcher–Goldfarb–Shanno (BFGS) algorithm was used in the geometry optimization process until the self-consistent-field (SCF) convergence condition was reached. The energy change per atom was set to 2×10⁻⁵ eV, the maximum force to 0.05 eV/Å, the maximum stress to 0.1 GPa, and the maximum displacement of atoms to 0.002 Å. To calculate the migration energy of the lithium ion, the nudge elastic band (NEB) [27] calculation was carried out using Transition State Theory, which is available in CASTEP.

RESULTS AND DISCUSSION

Structural and Mechanical Properties

LFP formed crystals with an orthorhombic structure and *Pnma* space groups (no. 62). In its primitive cell, it contained four formula units (Li₄Fe₄P₄O₁₆). The cell was formed by FeO₆ octahedra which were connected to PO₄ tetrahedra by edge-sharing. This created a one-dimensional lithium-ion channel in the (010) plane along the *y*-axis for it to intercalate and deintercalate throughout the cathode material. Figure 1 shows the crystal structure of LFP and Ca-LFP. The structure was constructed based on experimental data obtained from XRD refinement analysis with lattice parameters *a* = 10.349 Å, *b* = 6.031 Å, *c* = 4.722 Å, and volume, *V* = 294.718 Å³ [28]. The FP structure was created simply by removing the lithium ions in the cathode material.

To understand the effects of Ca substitution on structural properties, the lattice parameters of LFP before and after substitution were calculated. Table 1 shows the calculated lattice parameters of LFP and Ca-LFP in comparison with other theoretical and experimental data. For LFP and FP, the current work slightly overestimates the experimental lattice parameters. This is due to the effect of the Hubbard U parameter that tends to slightly overestimate lattice parameters and volume. The results are still acceptable and in agreement with previous theoretical studies [29]. Upon substituting the Fe atom with a Ca atom, the lattice parameters of LFP increased slightly. This is expected as the ionic radius of Ca²⁺ (1.00 Å) is bigger than that of Fe²⁺ (0.78 Å) [30]. Structural expansion provides a wider pathway for lithium diffusion, which should improve the charge/discharge rate.

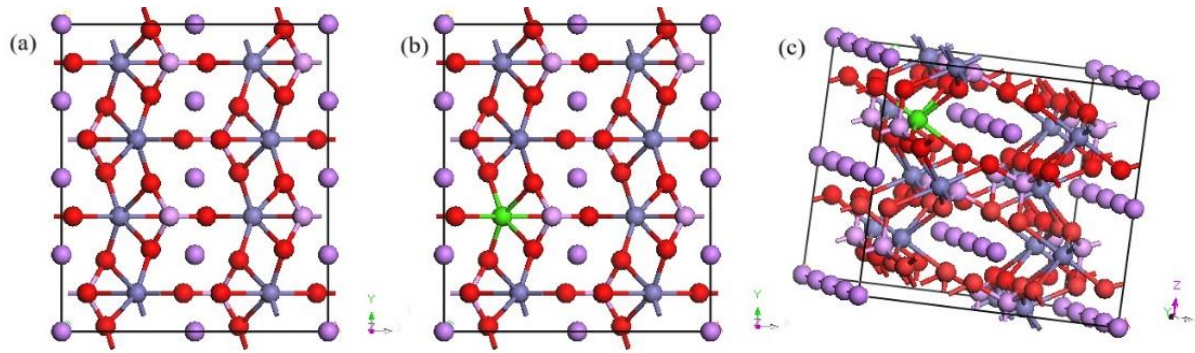


Figure 1. Crystal structures of (a) LFP, (b) Ca-LFP and (c) perspective view of Ca-LFP. The lithium ions, oxygen atoms, iron atoms, phosphorus atoms and calcium atoms are in purple, red, grey, pink and green, respectively.

Charging and discharging involves the extraction and insertion of lithium ions into the cathode material. During this process, the cathode undergoes contraction and expansion. This recurring process can degrade the battery and reduce its life cycle. Too much expansion can lead to instability and structural collapse. As lithium is extracted from the cathode material, the Fe²⁺ (0.78 Å ionic radius) is oxidized to Fe³⁺ (0.65 Å) causing shrinkage of the cathode material. This is one of the reasons why FP has a smaller volume compared to LFP.

The delithiation process reduced the calculated volume of Ca-LFP by 2.56 %. This is because the Ca²⁺ ion which replaced the Fe²⁺ does not oxidize during delithiation. As the amount of Fe³⁺ was reduced, the volume change was limited. With this behaviour, the life cycle of LFP should improve with Ca doping. However, there may be small reductions in theoretical capacity (160 mAhg⁻¹) as Ca²⁺ does not contribute to the redox reaction of Fe²⁺/Fe³⁺ which the battery depends on. Nonetheless, this reduction will be

compensated by the increase in electronic and ionic conductivity, which will be discussed later.

Formation energy, E_f , represents the energy change when a compound is formed from its basic elements in their standard state. Evaluating the formation energy is important to determine the thermodynamic stability of the doping system. The lower the formation energy, the more stable the system [32]. Formation energy is calculated based on Eq. (1), where E_{total} is the total energy of the doped system, E_{LFP} is the total energy of pristine LFP, and E_{Fe} and E_{Ca} are the chemical potentials of iron and calcium, respectively.

$$E_f = E_{total} - E_{LFP} + E_{Fe} - E_{Ca} \quad (1)$$

The formation energy for this doped system was -4.07 eV. Compared to previous studies which substituted different elements, this calculated value is acceptable [16,33,34]. Thus, this doped system is feasible and a thermodynamically stable system.

Table 1. Comparison of the lattice parameters of LFP and Ca-LFP with other theoretical and experimental data.

	a (Å)	$2b$ (Å)	$2c$ (Å)	V (Å ³)	ΔV (Å ³)	Ref
^a LiFePO ₄	10.33	12.02	9.39	1166.48	6.60 %	[31]
^a FePO ₄	9.82	11.58	9.58	1089.44		[31]
^b LiFePO ₄	10.44	12.14	9.49	1204.24	3.42 %	[29]
^b FePO ₄	9.986	11.82	9.762	1152.84		[29]
^b LiFePO ₄	10.44	12.15	9.48	1202.70	4.61 %	current work
^b FePO ₄	9.93	11.85	9.75	1147.28		current work
^b Ca-LiFePO ₄	10.45	12.16	9.48	1205.01	2.56 %	current work
^b Ca-FePO ₄	9.97	11.96	9.84	1174.12		current work

^a experimental, ^b theoretical

Table 2. Calculated elastic constants of LFP and Ca-LFP (in GPa).

	C ₁₁	C ₂₂	C ₃₃	C ₄₄	C ₅₅	C ₆₆	C ₁₂	C ₁₃	C ₂₃	Ref
LFP-cal.	138.9	198.0	173.0	36.8	50.6	47.6	72.8	52.5	45.8	[37]
LFP-cal.	149.1	187.0	178.2	39.1	50.2	46.7	71.4	58.9	48.0	[16]
LFP	140.0	182.4	169.5	37.8	47.2	48.8	67.3	53.4	44.2	current work
Ca-LFP	134.0	179.2	163.4	39.0	44.6	46.7	66.6	55.4	45.1	current work

The electrochemical performance of a cathode material relies on the stability of its material. Instability can lead to phase transformation and degradation. These can reduce the charge and discharge efficiency of the battery. To determine material stability, the elastic constant, C_{ij} is calculated. The elastic constant is commonly used to state the relationship between stress and strain in a material. Under the conditions of the elastic limit and small strains, C_{ij} serves as an indicator of the material's mechanical stability. In an orthotropic system, the elastic constants C_{11} , C_{22} , and C_{33} represent the linear compressive moduli along the X, Y, and Z direction respectively. C_{44} , C_{55} , and C_{66} correspond to the shear moduli in the (100), (010), and (001) planes, while C_{12} , C_{13} , and C_{23} are related to longitudinal and transverse contractions.

The calculated elastic constants C_{ij} for LFP and Ca-LFP are presented in Table 2. According to the Born criterion, an orthotropic crystalline system is mechanically stable if it satisfies the following conditions, [35, 36]:

$$\begin{aligned} C_{11} > 0, C_{22} > 0, C_{33} > 0, C_{44} > 0, C_{55} > 0, \\ C_{66} > 0, C_{11} + C_{22} + C_{33} + 2(C_{12} + C_{13} + C_{23}) \\ > 0, (C_{11} + C_{22} - 2C_{12}) > 0, (C_{11} + C_{33} - 2C_{13}) \\ > 0, (C_{22} + C_{33} - 2C_{23}) > 0 \end{aligned} \quad (2)$$

Based on the calculated lattice constants in Table 2, both LFP and Ca-LFP meet the stability conditions. This suggests that Ca doping is practical and does not reduce mechanical stability. To further evaluate the mechanical properties, two standard methods, the Voigt and Reuss methods, were employed to determine the bulk and shear moduli of LFP and Ca-LFP. The Voigt method assumes uniform strain, while the Reuss method assumes uniform stress. These two methods determine the upper and lower limits of the elastic modulus of the polycrystalline material. By averaging both values, the more accurate

expression of bulk modulus, B and shear modulus, G can be obtained.

$$B = \frac{B_V + B_R}{2} \quad (3)$$

$$G = \frac{G_V + G_R}{2} \quad (4)$$

The bulk modulus and shear modulus are then used to compute the Young's modulus, E , and Poisson's ratio, ν , using the following equations [35]:

$$E = \frac{9BG}{3B + G} \quad (5)$$

$$\nu = \frac{3B - 2G}{6B + 2G} \quad (6)$$

Table 3 provides a comparison of the mechanical properties of LFP and Ca-LFP. The bulk modulus reflects resistance to uniform compression, while the shear modulus indicates resistance to shear deformation, with higher values indicating greater material stiffness. Young's modulus highlights the material's resistance to deformation under stress. Prior to doping, there was a slight decrease in the mechanical properties (B , G , E) of LFP, signifying a small reduction in strength, stiffness, and ductility. In terms of Poisson's ratio, the Ca-LFP showed shear-induced deformation properties similar to LFP. The Pugh criterion, which is the ratio of B/G , was also used to assess material ductility and brittleness. A value above 1.75 indicates ductile behaviour, while a lower value indicates brittleness. For LFP and Ca-LFP, the B/G ratios were 1.91 and 1.95 respectively. This shows that Ca doping improved the ductility of the cathode material. Higher ductility indicates that the cathode material is less prone to cracking during prolonged cycling, which is important for its stability and efficiency.

Table 3. Shear Modulus (G), Bulk Modulus (B), Young Modulus (E), Poisson Ratio (ν) and Pugh Criterion, B/G values. G , B and E are in GPa.

	G_R	G_V	G	B_R	B_V	B	E	ν	B/G	Ref.
LFP	48.6	46.9	47.7	91.3	90.6	91.0	121.9	0.30	1.91	Present work
Ca-LFP	46.7	45.2	46.0	90.1	89.3	89.7	117.8	0.30	1.95	Present work
LFP-cal.	47.2	49.6	48.4	93	94.7	93.9	123.9	0.28	1.94	[37]
LFP-cal.	49.2	51.1	50.2	96.5	97.1	96.8	128.9	0.28	1.93	[32]

Microcracks or deformations may develop in materials due to crystal anisotropy. This behaviour has previously been discovered in LFP particles, and becomes more frequent with prolonged cycling [38]. One of the factors that might contribute to these defects is elastic anisotropy. To calculate this effect, Ranganathan et al. [39] introduced a universal elastic anisotropy index (A^U):

$$A^U = 5 \frac{G_V}{G_R} + \frac{B_V}{B_R} - 6 \quad (5)$$

where G_V , G_R , B_V and B_R represent the calculated values from the Voigt and Reuss methods. Isotropy is present in a crystal if $A^U = 0$. If the value is greater than 0, the degree of anisotropy will increase in the crystal structure. In the present study (Table 4), the anisotropy index of LFP was 0.19, in agreement with previous work [32]. After Ca substitution, the index reduced to 0.17, indicating an increase in the degree of isotropy. Another anisotropy factor which has been defined by Chung and Buessem [40] is as follows:

$$A_B = \frac{B_V - B_R}{B_V + B_R} \quad (6)$$

$$A_G = \frac{G_V - G_R}{G_V + G_R} \quad (7)$$

The more the value approaches zero, the more isotropic the material. The values obtained after doping suggest that the material's isotropy remained

Unchanged (Table 4). The following are shear anisotropy factors by Ravindran.

$$A_1 = 4C_{44}/(C_{11} + C_{33} - 2C_{13}) \quad (8)$$

$$A_2 = 4C_{55}/(C_{22} + C_{33} - 2C_{23}) \quad (9)$$

$$A_3 = 4C_{66}/(C_{11} + C_{22} - 2C_{12}) \quad (10)$$

These provide a measure of the degree of anisotropy in the bonding between atoms in different planes [41]. A_1 is for the (100) shear plane in the $\langle 010 \rangle$ and $\langle 011 \rangle$ directions, A_2 for the (010) shear plane in the $\langle 001 \rangle$ and $\langle 101 \rangle$ directions, and A_3 for the (001) shear plane in the $\langle 010 \rangle$ and $\langle 110 \rangle$ directions. If the values of A_1 , A_2 , and A_3 are equal to 1, the material will exhibit isotropy, and the more the value deviates from 1, the higher the degree of anisotropy. As can be seen in Table 4, prior to doping with Ca, LFP exhibited some degree of anisotropy in the (100) and (010) planes and tended to be isotropic in the (001) plane. After Ca doping, there was an improvement in the degree of isotropy in the (100) plane. In the (010) and (001) planes, no significant change was observed. Overall, Ca doping positively improved the degree of isotropy in LFP which lowers the likelihood of microcrack formation. It must be noted that a tiny crack can negatively degrade the electrochemical performance of the cathode material. Previous work has shown that microcracks caused some LFP particles to be electronically disconnected and suffer a loss in capacity [42]. Thus, improving the degree of isotropy should help to improve the life cycle of the cathode material.

Table 4. Shear Anisotropy Factors (A_1 , A_2 and A_3), Elastic Anisotropy (A_B and A_G) and Universal Elastic Anisotropy index (A_U).

	LFP	Ca-doped	LFP [16]	LFP [32]	LFP [37]
A_1	0.75	0.84	0.74	0.73	0.712
A_2	0.61	0.60	0.74	0.77	0.724
A_3	1.04	1.04	0.96	1.06	0.995
A_B	0.00	0.00	-	0.003	-
A_G	0.02	0.02	-	0.0019	-
A_U	0.19	0.17	-	0.20	-

Electronic Properties

As mentioned in a previous work by Zhang et al., the major factor for determining the electronic conductivity of a solid is the band gap [43]. This energy gap can be obtained from the density of state calculation. Figure 2 displays the density of state of LFP and FP. In between the conduction bands and valence bands of LFP and FP, there is an energy gap of 3.696 eV and 1.447 eV, respectively. These values are in agreement with previous theoretical and UV-Vis experimental data [34, 44, 45]. The calculated results were generated using a Hubbard U parameter with a value of 4.3 eV. Without it, the calculated energy gap would be underestimated by about 0.5 eV for LFP, which would incorrectly indicate this cathode material as metallic. Incorporation of Hubbard U overcomes the incomplete cancellation of the electronic self-interaction error of the delocalized 3d state. The energy gap of 3.696 eV can be considered as a wide band gap, impeding the self-

generation of electrons or holes. This drawback gives LFP insulating behaviour leading to poor electronic conductivity, which is the main problem with LFP cathode materials. To solve this issue, foreign impurities may be introduced into the cathode material to tune its conductivity.

Figure 3 displays the calculated density of state of Ca-LFP and Ca-FP. The presence of Ca as a dopant altered the electronic structure of LFP. There was a small contribution from Ca state to the density of state, particularly at the conduction band. The Ca state weakly overlapped with the Fe state and seemed to play a role in reducing the energy gap of LFP and FP to 2.164 eV and 1.110 eV, respectively. As can be seen in the Figure 2 (a), with the involvement of the Ca state, the Fe state in LFP, a prominent peak at ~5 eV shifted to ~2.5 eV (Figure 3 (a)). This reduced gap allows electrons to easily transition from the valence band to the conduction band, as less energy is required [9, 16, 19].

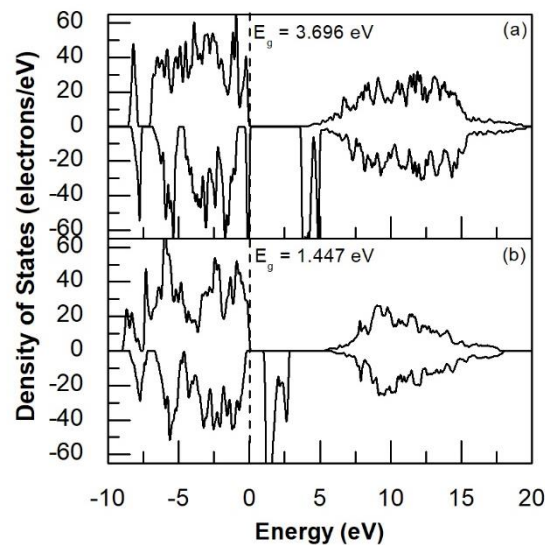


Figure 2. The density of state (DOS) of (a) LFP and (b) FP.

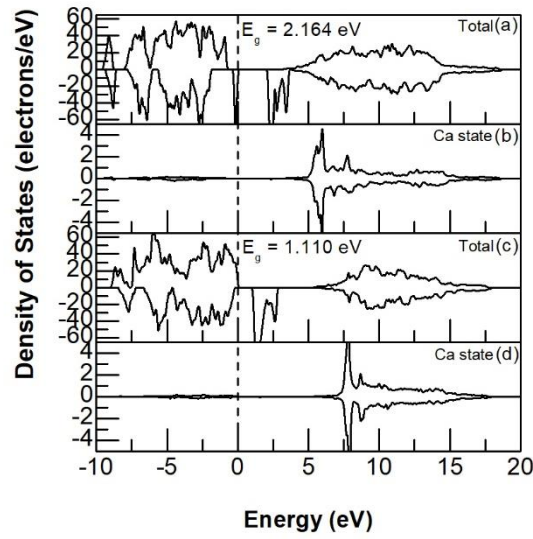


Figure 3. The calculated DOS of Ca-LFP and Ca-FP: (a) The DOS of Ca-LFP, (b) the Ca atom state in the LFP DOS, (c) the DOS of Ca-FP and (d) Ca atom state in the FP DOS.

The reduction in the energy gap implies higher electrical conductivity. This is in agreement with previous work. Liu et al. [22] performed Ca doping on LiFe_{0.5}Mn_{0.5}PO₄ and found that it greatly increased electronic conductivity by reducing charge transfer resistance, and increased lithium-ion diffusion. Similar trends were also found by Fey et al. [26] in their work on a Ca doped LFP which showed an increase in electronic conductivity and lithium-ion diffusion. This confirms that Ca doping effectively improved the electrochemical performance of LFP.

It is worth noting that although the electronic conduction in LiFePO₄ is dominated by thermally activated polaron hopping rather than intrinsic band conduction [46], the calculated energy gap, E_g is often used as a qualitative descriptor of electronic conductivity. In intrinsic semiconductors, the carrier concentration, and thus the conductivity, σ follow an Arrhenius-type dependence, $\sigma \propto \exp\left(-\frac{E_a}{kT}\right)$, where $E_a \approx E_g/2$. This means that a narrower band gap corresponds to higher conductivity. While this relationship does not strictly apply to polaronic materials, dopant-induced band gap narrowing in LiFePO₄, generally indicates the formation of intermediate states closer to Fermi level, which reduces the energy gap, E_g and effective activation energy, E_a for polaron hopping and thereby enhances electronic conductivity. This has been demonstrated in previous work by Chung et al. [47] on Mg, Zr and Nb doping, which showed that the activation energy was reduced

to ~60-80 meV compared to undoped LiFePO₄ (500 meV), and this was accompanied by an improvement in conductivity.

Lithium Migration

In addition to electronic conductivity, ionic conductivity also plays an important role in the electrochemical performance of the cathode material, especially in terms of rate capability. This is related to Li ion mobility inside the cathode material which can be probed via the migration energy, E_m calculated based on transition state. Figure 4 shows the comparison of the migration energy of LFP and Ca-LFP along the [010] direction. This direction is the most stable pathway and is responsible for lithium ion diffusion in the LFP cathode material [48]. The migration energy path looks arch-like due to electrostatic interactions between Li and its neighbouring Fe or Ca ion. Prior to doping the migration energy of the lithium ion was 0.532 eV. After Ca doping, the migration energy reduced to 0.459 eV. This value can be used to calculate the diffusion coefficient, D using the Arrhenius equation below:

$$D = a^2 \nu \exp\left(-\frac{E_m}{k_B T}\right) \quad (11)$$

where a is the jump distance along the b-axis (3.0 Å), ν is attempt frequency (10^{13} Hz), k_B is the Boltzmann constant, T is the temperature of the system (300 K), and E_m is the migration energy.

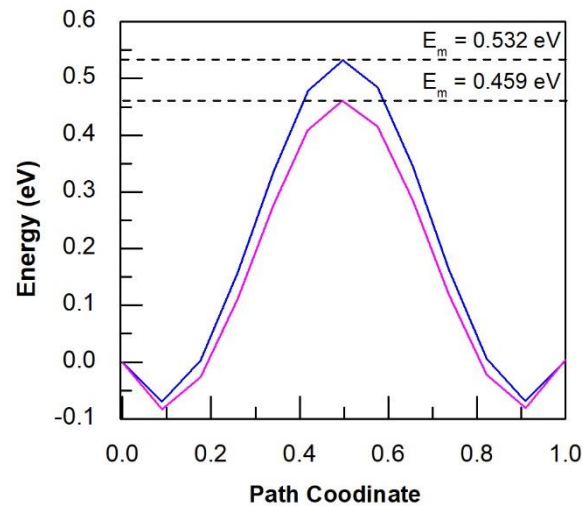


Figure 4. The migration energy of the Li ions in LFP (blue line) and Ca-LFP (pink line).

Table 5. Calculated Diffusion Coefficients for LFP and Ca-LFP using the Arrhenius Equation.

	Migration Energy, E_m (eV)	Diffusion Coefficient, D (cm ² s ⁻¹)
LFP	0.532	1.04×10^{-11}
Ca-LFP	0.459	1.75×10^{-10}

*Parameters use: $a = 3.0 \text{ \AA}$, $\nu = 10^{13} \text{ Hz}$ and $T = 300 \text{ K}$

For LFP and Ca-LFP the calculated diffusion coefficients (Table 5) were $1.04 \times 10^{-11} \text{ cm}^2 \text{ s}^{-1}$ and $1.75 \times 10^{-10} \text{ cm}^2 \text{ s}^{-1}$, respectively. This indicates that Ca doping improved the ionic conductivity of the lithium ion. This has been shown in a previous study where LiFeMnPO₄ demonstrated higher diffusion coefficients compared to LiFePO₄ during the discharge process [22]. Zhang et al. also demonstrated that Ca doping in LiFePO₄ significantly accelerated lithium migration due to broadening of the lithium-ion migration pathway [49]. Ca doping increased the lattice parameters and volume of LFP and provided an easy pathway for lithium migration. This is known as the “pillar effect” which reduces energy barriers for lithium ion diffusion, making it easier for them to move within the material [50], thus increasing the diffusion rate and enhancing the rate performance of the cathode.

CONCLUSION

The effects of Ca doping on the LFP cathode material were investigated. It was found that Ca doping improved mechanical and electrochemical performance. Ca-LFP satisfied the thermodynamic stability requirements, which indicated that this doping system was feasible. Based on the stability conditions, Ca doping did not compromise mechanical stability. Even with a slight

decrease in mechanical properties (B , G , E), there was an improvement in ductility and the degree of isotropy, which should prevent microcrack formation after prolonged usage. In terms of electronic properties, the energy gap values for LFP and FP were reduced to 2.164 eV and 1.110 eV, indicating improved electronic conductivity. Moreover, based on transition state theory, the migration energy of Ca-doped LFP reduced from 0.532 eV to 0.459 eV. This contributed to an increase in the diffusion coefficient from $1.04 \times 10^{-11} \text{ cm}^2 \text{ s}^{-1}$ to $1.75 \times 10^{-10} \text{ cm}^2 \text{ s}^{-1}$, indicating faster lithium ion diffusion, which should offer a higher rate capability. To further confirm these results, experimental synthesis and electrochemical testing should be performed in the future. Other explorations such as a defect study, and doping optimization should also be attempted to seek deeper insights into the practical implementation of a doping strategy for next generation energy storage devices.

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